## Chapt2: Configuration of the Atom: Rutherford's Model

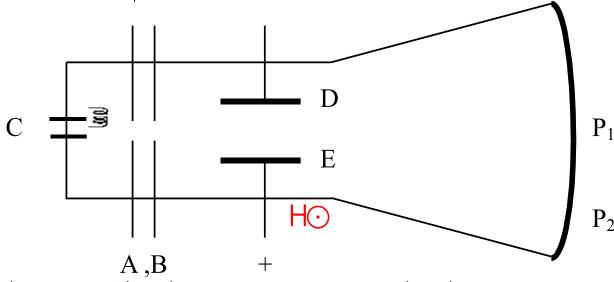
- 2. 1 Background
- 2.2 Emergence of the Rutherford's Model
- 2.3 Rutherford scattering Formula
- 2.4 Experimental test of the RF Formula.
- 2.5 Summary of the Nuclear planet model

#### 2-1 Background

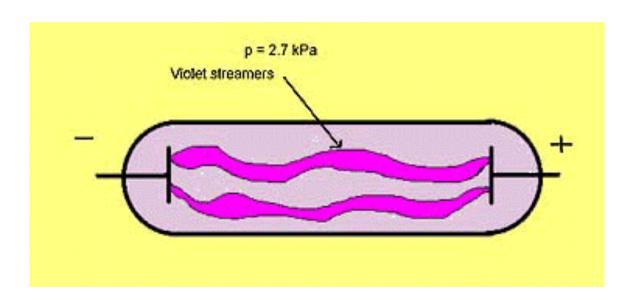
- 1. The discovery of the electron
- 2. The charge and mass of the electron
- 3. The size of Atom

#### 2.1.1 Discovery of the electron

• Thomson's discharge tube



- Cathode rays (C)  $\Rightarrow$  Slits (A, B)  $\Rightarrow$  paralle plates (D, E)  $\Rightarrow$  screen
- D, E with E  $\Rightarrow$  rays  $P_1 \rightarrow P_2 \Rightarrow$  negative charged
- With H  $\Rightarrow$  rays  $P_2 \rightarrow P_1 \Rightarrow Hev = Ee \Rightarrow v = E/H$
- Without  $E \Rightarrow \text{rays}$  radius  $r \Rightarrow mv^2/r = Hev \Rightarrow e/m = v/Hr$



#### Thomson's Gas discharge tube

- In the late 1800s several experiments were performed on electrical discharges in gases at low pressure.
- It was established that a stream of some kind of "rays" called cathode-rays was emitted by any electrode at high negative potential in a vacuum tube.
  - What were these rays?

#### 2.1.2 Charge & Mass of Electron

- 0il-drop experiment by Millikan:
  - Measured charge :  $e = 1.6 \times 10^{-19} \text{ C}$  $\Rightarrow$ electron's mass  $m_e = 9.1 \times 10^{-31} \text{ kg}$
  - Millikan found charge is quantized
    - Any charge can only be an integral multiple of *e* (smalled charge)

given 
$$\frac{e}{m_p} \implies \frac{m_e}{m_p} = \frac{1}{1836}$$

- 0il drop experiment (2)
  - -Atom is neutral, with negative charged electrons, ⇒ there must be postive charged matter ( proton)

Atom = Postive + Negative + Neutral

#### 2. 1. 3 Avogadro's Constant

- $N_{\Lambda}$  gives the number of atoms in one Mole
- Atomic mass unit [u]
  - $-1[u] \equiv$  The mass of one <sup>12</sup>C /12

$$= \frac{12\bar{R}}{12 \times N_A} = \frac{1}{N_A} [\bar{R}] = 1.66 \times 10^{-24} [\bar{R}]$$

- Atomic mass  $M_A$  [u] = atomic quantity [u] = Mass number A [u]
- $-1g=uN_A$ ;  $F=eN_A$   $R=kN_A$
- Macro -- N<sub>A</sub> --- Micro

#### 2.1.4 The size of atoms

• Atomic radius

- Atomic volum = 
$$\frac{4}{3}\pi r^3$$

= Mass of an atom/atomic density =  $\frac{A/N_A}{\rho}$ 

$$\Rightarrow$$
 radius  $r = \left(\frac{3A}{4\pi\rho N_A}\right)^{\frac{1}{3}}$ 

order of magnitude:  $r \sim 10^{-10} \text{ m} = 1 \text{ Å}$ 

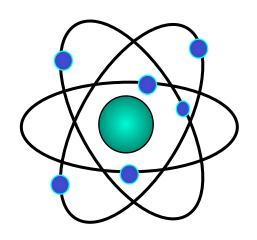
- 2.2 Emergence of Rutherford's Model
  - 1. Atomic charges' distribution?
  - 2.  $\alpha$  scattering
  - 3. Scattering experiments

## 2.2.1 How are charges distributed in atoms

• Thomas atomic model postive charges are distributed uniformly over the whole atom, electrons are embedded

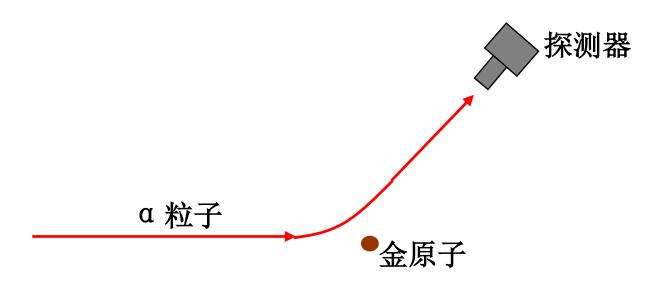
#### Rutherford model

Atom is composed of by nuclei and circling around electrons



#### $2.2.2 \alpha$ scatting experiment

α 粒子散射实验:



- results:
  - Most  $\alpha$  particles have scattering angle:  $2^{\circ} 3^{\circ}$
  - few of  $1/8000 \alpha$  with SA: >  $90^{\circ}$
- puzzled: like a gun bullet was bounced back by a paper

#### 2.2.3 Quantitative explaination

- Thomson's model
  - Force of Positive charge Ze of atom on the α particle

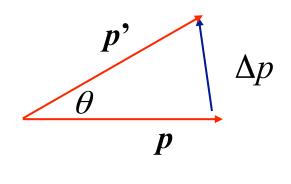
$$F = \begin{cases} \frac{1}{4\pi\varepsilon_0} \frac{2Ze^2}{r^2} & r \ge R \\ \frac{1}{4\pi\varepsilon_0} \frac{2Ze^2}{R^3} r & r < R \end{cases}$$
Radius of atom

#### The Max. Force of Positive charge Ze of atom on the $\alpha$ particle

$$F = \frac{1}{4\pi\varepsilon_0} \frac{2Ze^2}{R^2}$$

-Scat. Angle

$$\theta = \frac{\Delta p}{p}$$



Momentum change~ F t

(2R/v)

#### So we have

$$\frac{\Delta p}{p} = \frac{2FR/v}{m_{\alpha}v} = \frac{2Ze^{2}/(4\pi\varepsilon_{0}R)}{\frac{1}{2}m_{\alpha}v^{2}} E_{\alpha}$$

$$\approx \frac{2Z \times 1.44 \text{fmMeV}/0.1 \text{nm}}{E_{\alpha}(\text{MeV})} \approx 3 \times 10^{-5} \frac{Z}{E_{\alpha}} \text{rad}$$

#### 2.2.3 Electron's effect

-Electron and a particle's (head on )

#### P and E conservation $\rightarrow$

$$m_{\alpha}v_{0} = m_{\alpha}v_{1} + m_{e}v_{e}$$
,  $\frac{1}{2}m_{\alpha}v_{0}^{2} = \frac{1}{2}m_{\alpha}v_{1}^{2} + \frac{1}{2}m_{e}v_{e}^{2}$ 

In coming  $\alpha$ 

Outgoing  $\alpha$ 

**Outgoing electron** 

$$\Rightarrow v_e = \frac{2m_\alpha v_0}{m_\alpha + m_e} \approx \frac{2m_\alpha v_0}{m_\alpha} = 2v_0$$

$$\Rightarrow \Delta p = m_\alpha v_0 - m_\alpha v_1 = m_e v_e \approx 2m_e v_0$$

#### therefore

- Momentum change

$$\Delta p = m_{\alpha} v_0 - m_{\alpha} v_1 = m_e v_e \approx 2m_e v_0$$

$$p = m_{\alpha} v_0$$

$$\Rightarrow \frac{\Delta p}{p} \approx \frac{2m_e v_0}{m_{\alpha} v_0} = \frac{2m_e}{m_{\alpha}} \approx \frac{1}{4000} \approx 10^{-4}$$

Deflection angle of  $\alpha$  is very tiny

Almost no contribution

-  $\alpha$ -particle on Gold atom

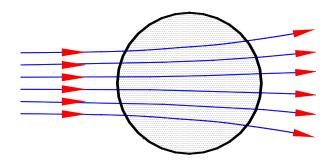
$$E_{c} = 5 \text{MeV}$$
 **Z**=**79**

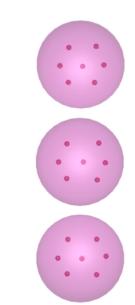
$$\theta = \frac{\Delta p}{p} \approx 3 \times 10^{-5} \frac{Z}{E_{\alpha}} \text{ rad} < 10^{-4} \frac{Z}{E_{\alpha}} \text{ rad} < 10^{-3} \text{ rad}$$

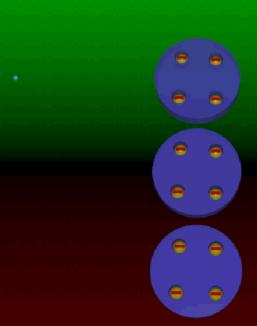
- One scattering angle  $10^{-3}$  rad
- -Multiple scatterings do not result large angle
  - Random in direction
- Thomson model contradicts with experiment result!

#### Thomson model predicts:

- SA 3°probabilty less than 1%
- SA greater than 90°, about 10<sup>-3500</sup>。







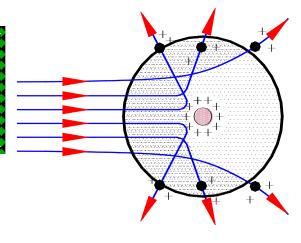
#### 1/8000

R

- most with small SA, about 1/8000 with SA>90°;
- Very few with 180



Positive charges locate on the central of atom



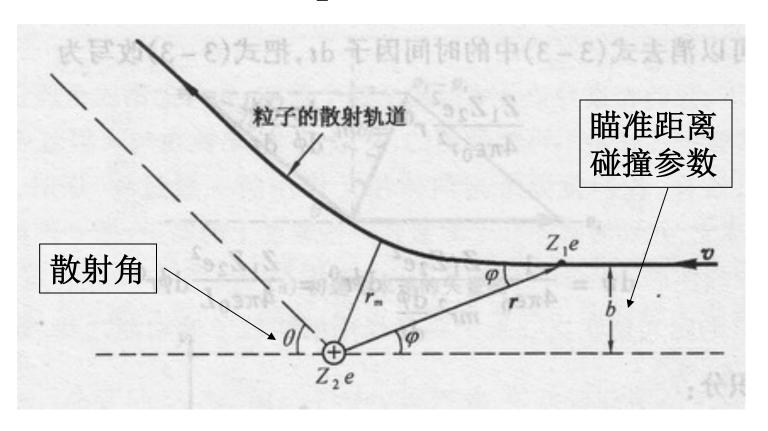
2.3 Rutherford Scattering Formula

1. Derivation of Coulomb formula

2. Derivation of Rutherford formula

#### 2-3-1 Coulomb scattering formula

• Scattering of particle with  $E, Z_1e$  by a target nucleus of charge  $Z_2e$ 



Coulomb scattering formula

$$b = \frac{a}{2}ctg\frac{\theta}{2}$$
 Impact parameter

$$a = \frac{Z_1 Z_2 e^2}{4\pi \varepsilon_0 E}$$

Scattering factor

#### • Assumptions:

- 1. Single scattering
- 2. Point charge and only Coulomb force
- 3. Extranuclear elctrons may be negnected
- 4. Target nucleus is rest

#### • According to

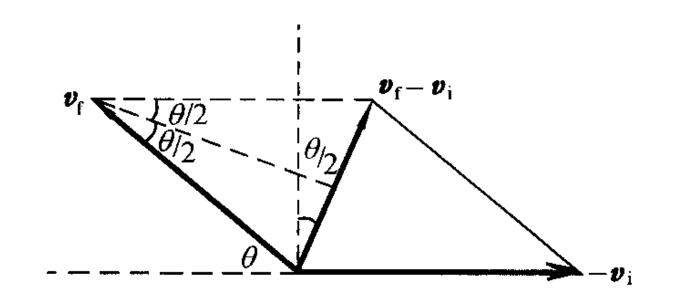
$$F = ma^{V} \rightarrow \frac{Z_1 Z_2 e^2}{4\pi\varepsilon_0 r^2} V_0 = m \frac{dv}{dt}$$

$$mr^2 \frac{\mathrm{d}\phi}{\mathrm{d}t} = L \leftarrow \Box$$
 Central force  $\rightarrow$  L conservation

$$\frac{Z_1 Z_2 e^2}{4\pi\varepsilon_0 r^2} \overset{\mathbf{v}_0}{r} = m \frac{\mathrm{d}^{\mathbf{V}}_{\mathbf{v}}}{\mathrm{d}\boldsymbol{\phi}} \frac{\mathrm{d}\boldsymbol{\phi}}{\mathrm{d}t}$$

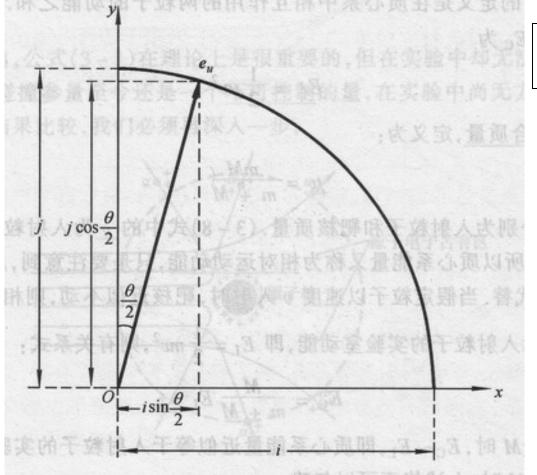
$$dv = \frac{1}{4\pi\varepsilon_0} \frac{Z_1 Z_2 e^2}{mr^2} d\phi r^{V_0} = \frac{Z_1 Z_2 e^2}{4\pi\varepsilon_0 L} d\phi r^{V_0}$$

$$\int d\mathbf{v} = \frac{Z_1 Z_2 e^2}{4\pi\varepsilon_0 L} \int_r^{\mathbf{v}_0} d\phi; \qquad \int d\mathbf{v} = \mathbf{v}_f - \mathbf{v}_i = \begin{vmatrix} \mathbf{v} & \mathbf{v} \\ \mathbf{v}_f - \mathbf{v}_i \end{vmatrix} \frac{\mathbf{v}}{\mathbf{e}_u}$$



$$E = \frac{1}{2}m|\mathbf{v}_{i}|^{2} = \frac{1}{2}m|\mathbf{v}_{f}|^{2} \rightarrow |\mathbf{v}_{i}| = |\mathbf{v}_{f}| = |\mathbf{v}_{f}| = \mathbf{v}; \quad |\mathbf{v}_{f} - \mathbf{v}_{i}| = 2v\sin\frac{\theta}{2}$$

$$\int_{0}^{\mathbf{V}_{0}} d\phi = \int_{0}^{\pi - \theta} (i \cos \phi + i \sin \phi) d\phi = 2 \cos \frac{\theta}{2} \left( i \sin \frac{\theta}{2} + i \cos \frac{\theta}{2} \right)$$



$$v\sin\frac{\theta}{2} = \frac{1}{4\pi\varepsilon_0} \frac{Z_1 Z_2 e^2}{L} \cos\frac{\theta}{2} = \frac{1}{4\pi\varepsilon_0} \frac{Z_1 Z_2 e^2}{mvb} \cos\frac{\theta}{2}$$

$$\rightarrow$$

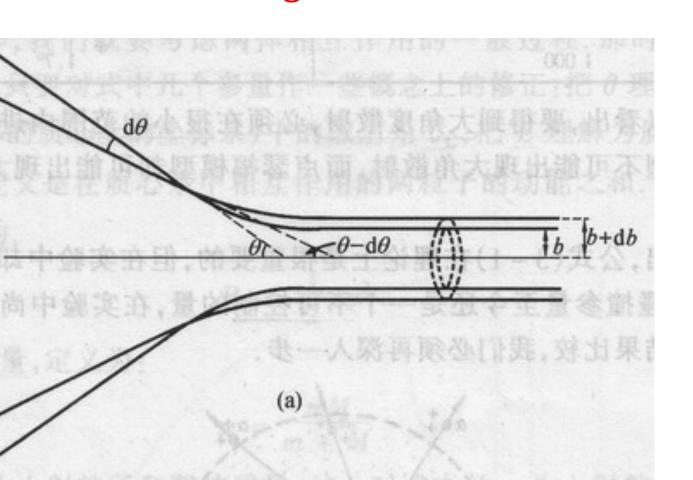
$$b = \frac{a}{2} \cot \frac{\theta}{2} \qquad a = \frac{Z_1 Z_2 e^2}{4\pi \varepsilon_0 E} \qquad E = \frac{mv^2}{2}$$

$$b \sim \theta$$
; b \rightarrow  $\theta \downarrow$ ; b \lor  $\theta \uparrow$ 

#### 2-3-2 Rutherford's formula (1)

 $\alpha$  particle :  $b \sim b + db \rightarrow \theta \sim \theta - d\theta$ 

 $\alpha:b\sim b+db \text{ ring area} \rightarrow \theta \sim \theta-d\theta \text{ hollow cone}$ 



prob. on the
ring of

- Target foil with thickness t,A, number density n, mass density  $\rho$   $2\pi b \mid db \mid$
- Area of the ring
- $\alpha$  particle on the ring

$$\frac{2\pi b |db|}{A} = \frac{2\pi}{A} \frac{a}{2} \cot \frac{\theta}{2} \left| -\frac{a}{2} \csc^2 \frac{\theta}{2} \frac{1}{2} d\theta \right|$$

$$= \frac{2\pi a^2 \sin\theta d\theta}{16A \sin^4 \frac{\theta}{2}}$$

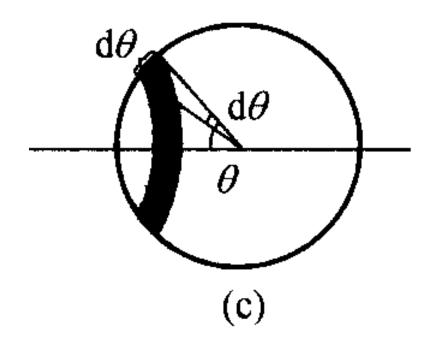
#### • Hellow cone cubic angle $\sim d\theta$

$$ds = 2\pi r \sin\theta \ rd\theta;$$

$$d\Omega = \frac{ds}{r^2} = 2\pi \sin\theta d\theta$$



$$\frac{2\pi b |db|}{A} = \frac{a^2 d\Omega}{16A\sin^4 \frac{\theta}{2}}$$



- Many rings in a foil: nucleus ~ring;
- volum: At; number of rings: Atn
- $\alpha$  on Atn, with the same  $\theta$
- The prob. Of one  $\alpha$  hiting on the foil with  $\theta \sim \theta d\theta$

$$dp(\theta) = \frac{a^2 d\Omega}{16A \sin^4 \frac{\theta}{2}} nAt$$

• N  $\alpha$  scattering on a foil  $\theta \sim \theta - d\theta$ 

$$dN' = N \frac{a^2 d\Omega}{16A \sin^4 \frac{\theta}{2}} nAt = ntN \left( \frac{1}{4\pi\varepsilon_0} \frac{Z_1 Z_2 e^2}{4E} \right)^2 \frac{d\Omega}{\sin^4 \frac{\theta}{2}}$$

• Differential Cross-section

$$\sigma_{C}(\theta) = \frac{d\sigma(\theta)}{d\Omega} = \frac{dN'}{Nntd\Omega} = \left(\frac{1}{4\pi\varepsilon_{0}} \frac{Z_{1}Z_{2}e^{2}}{4E}\right)^{2} \frac{1}{\sin^{4}\frac{\theta}{2}}$$

The Ns of scattered particles per number of atoms, per unit area of the target per incident particle per solid angle

• Understanding the D-cross section

$$\sigma_{C}(\theta) = \frac{d\sigma(\theta)}{d\Omega} = \frac{dN'}{Nntd\Omega} = \left(\frac{1}{4\pi\varepsilon_{0}} \frac{Z_{1}Z_{2}e^{2}}{4E}\right)^{2} \frac{1}{\sin^{4}\frac{\theta}{2}}$$

- For per unite area of the target nucleus, each incident particle, each unite cubic angle, the number of scatterings
- -Effective scattering area of every atom into the solid angle element in  $\theta$  direction
- Dimension : area  $(m^2/sr)$
- Cross-section: **b**):  $1b = 10^{-28}m^2, 1mb = 10^{-31}m^2$

#### 2-4 Experimental verification

- 1. Geiger-Marsden Experiment
- 2. Nuclear size

### Geiger-Marsden Experiment

$$\sigma_{C}(\theta) = \frac{d\sigma(\theta)}{d\Omega} = \frac{dN'}{Nntd\Omega} = \left(\frac{1}{4\pi\varepsilon_{0}} \frac{Z_{1}Z_{2}e^{2}}{4E}\right)^{2} \frac{1}{\sin^{4}\frac{\theta}{2}}$$
• fixing (\alpha source, target) 
$$\frac{dN'}{d\Omega} \propto \frac{1}{\sin^{4}\frac{\theta}{2}}$$

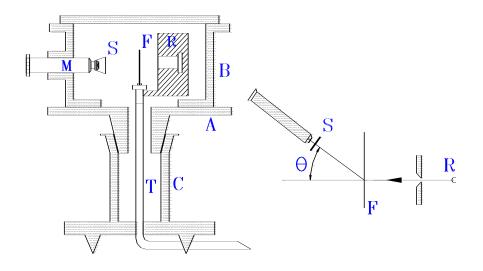
• Fixing (
$$\alpha$$
 source, material of target, angle)  $\frac{dN'}{d\Omega} \propto t$ 

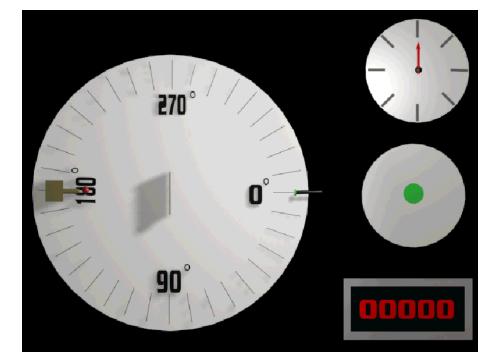
• same(target, angle) 
$$\frac{dN'}{dQ} \propto \frac{1}{F^2}$$

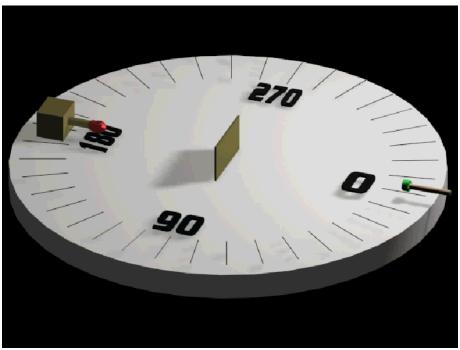
• same( $\alpha$  source, angle, nt)  $\frac{dN'}{d\Omega} \sim Z_2^2$ 

# 实验装置和模拟实验

- R
- **F**
- **S**
- **B**
- A
- C
- T
- **M**







• The experiment verified above relations!

#### 2-4-2 Nuclear size estimation (1)

$$\begin{cases} \frac{1}{2} m v_0^2 = E = \frac{1}{2} m v_m^2 + \frac{Z_1 Z_2 e^2}{4\pi \varepsilon_0 r_m} \\ m v_0 b = m v_m r_m = L \end{cases}$$

$$r_m^2 = \frac{L^2}{2mE} + \frac{Z_1 Z_2 e^2}{4\pi\varepsilon_0 E} r_m; \quad r_m^2 - ar_m - b^2 = 0$$

$$a = \frac{Z_1 Z_2 e^2}{4\pi \varepsilon_0 E}, \qquad \frac{m^2 v_0^2 b^2}{2m \cdot \frac{1}{2} m v_0^2} = b^2$$

#### 2.4.2 Estimation of the Size(2)

• solve

$$r_{m}^{2} - ar_{m} - b^{2} = 0$$

$$r_{m} = \frac{a \pm \sqrt{a^{2} + 4b^{2}}}{2}, r_{m} > 0 \Rightarrow$$

$$r_{m} = \frac{a}{2} + \frac{a}{2}\sqrt{1 + \frac{4b^{2}}{a^{2}}} = \frac{a}{2}(1 + \sqrt{1 + \cot^{2}\frac{\theta}{2}})$$

$$= \frac{a}{2}(1 + \frac{1}{\sin\frac{\theta}{2}})$$

$$b = \frac{a}{2}\cot\frac{\theta}{2}$$

$$\theta = 180^{\circ} \to r_m = a : 1.20 \text{fm}$$

#### 2.5 On Nuclear model

- Significances:
  - Nuclear struncture of atom
  - New method black body method
- difficulties:
  - Stable (accelerating→radiation)
  - Identical (atom ←> solar system, IC)
  - Regeneration